

The darkness crumbles away.
It is the same old druid time as ever,
Only a live thing keeps my hand,
A queer aridonic rat,
Chapter 6: Oscillations

As I pulled the puppet's pappy.
To stick behind my ear.
To roll out, they would shoot you if they knew
Your cosmopolitan sympathies.
Now you have touched this English hand.

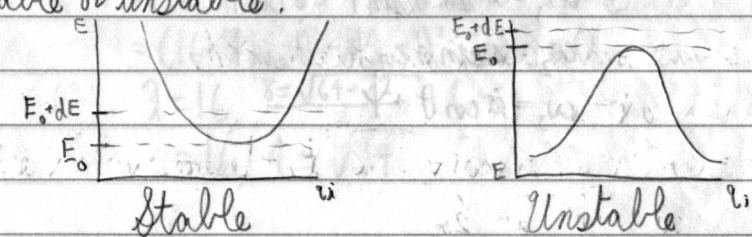
Section 1. Formulation of the Problem

Imagine we have a conservative system whose potential energy is solely a function of position. The system is said to be at equilibrium if

$$Q_i = -\left(\frac{\partial V}{\partial q_i}\right)_0 = 0 \quad (6.1)$$

We can imagine that if the system starts at equilibrium, it will stay at equilibrium (if you buy a spring and set it on a table, will it spontaneously start to move?). However, what will happen if perturb the system slightly?

If we plot the potential, we can determine if the motion is stable or unstable.



We can imagine a particle at equilibrium (E_0) is pushed a little to the right. If the motion is stable, it will return to equilibrium. If the motion is unstable, the particle will fly off.

For small deviation, we can expand the potential as a Taylor series about the equilibrium

$$V(q_1, \dots, q_n) = V(q_{01}, \dots, q_{0n}) + \left(\frac{\partial V}{\partial q_i}\right)_0 \eta_i + \frac{1}{2} \left(\frac{\partial^2 V}{\partial q_i \partial q_j}\right)_0 \eta_i \eta_j + \dots \quad (6.3)$$

$$\approx \frac{1}{2} \left(\frac{\partial^2 V}{\partial q_i \partial q_j}\right)_0 \eta_i \eta_j = \frac{1}{2} V_{ij} \eta_i \eta_j \quad (6.4)$$

since we set the equilibrium position to 0 and from the condition set in (6.1).

We can similarly treat the kinetic energy

$$T = \frac{1}{2} T_{ij} \dot{\eta}_i \dot{\eta}_j \quad (6.6)$$

$$L = \frac{1}{2} T_{ij} \dot{\eta}_i \dot{\eta}_j - \frac{1}{2} V_{ij} \eta_i \eta_j \quad (6.7)$$

$$= \frac{1}{2} (T_{ij} \dot{\eta}_i^2 - V_{ij} \eta_i \eta_j) \quad (6.9)$$

You will do the same to a German.
Soon, no doubt, if it be your pleasure
To cross the sleeping green between.
It seems as if you inwardly grin as you pass
Strong eyes, fine limbs, haughty athlete,

Less chance than you for life,
Bonds to the whims of murder,
Sprawled in the bowels of the earth,
The torn fields of France.
What do you see in our eyes

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\eta}_i} \right) - \frac{\partial L}{\partial \eta_i} = 0$$

$$\frac{d}{dt} (T_{ij} \dot{\eta}_j) - (-V_{ij} \eta_j) = 0 \quad (6.10)$$

Section 2. The Eigenvalue Equation and the Principal Axis Transformation

The solution to the harmonic oscillator is a known quantity, and one we've used in the past. Say we have a spring whose Lagrangian can be written as

$$L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$$

$$m \ddot{x} + kx = 0$$

$$x(t) = A \exp(i \omega t) + B \exp(-i \omega t)$$

$$\omega = \sqrt{k/m}$$

Now let's say we had some damping force $\vec{F} = -\gamma \dot{x}$

$$m \ddot{x} + \gamma \dot{x} + kx = 0$$

$$x(t) = \exp(-\beta t) [A \exp(\sqrt{\beta^2 - \omega_0^2} t) + B \exp(-\sqrt{\beta^2 - \omega_0^2} t)]$$

$$\omega = \sqrt{k/m}, \beta = \gamma/2m$$

There are three interesting cases for damped harmonic motion

Underdamped: $\beta^2 < \omega_0^2$

$$x(t) = C \exp(-\beta t) \cos(\sqrt{\omega_0^2 - \beta^2} t + \phi)$$

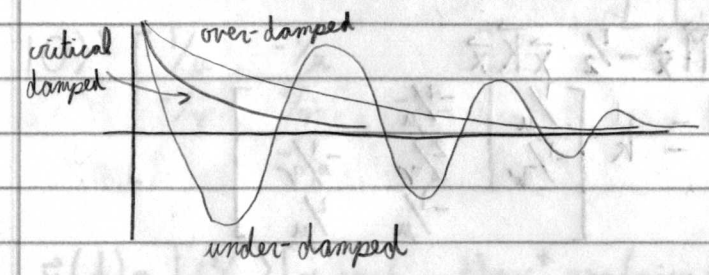
$$x(\infty) = 0$$

And we can convince ourselves that energy is not conserved because of the dissipative forces.

Overdamped: $\beta^2 > \omega_0^2$

Critical damping: $\beta^2 = \omega_0^2$

$$x(t) = C \exp(-\beta t)$$



At the striking iron and flame
 Hurl'd through still heavens?
 What quiver - what heart aght!
 Poppies, whose roots are in man's veins
 Drop, and are ever dropping;

But mine in my ear is safe -
 Just a little white with the dust.
 - Isaac Rosenberg (Break of Day in
 the Trenches, 1916)

If we have some external driving force

$$m\ddot{x} + \gamma\dot{x} + kx = F(t)$$

Section 3 Frequencies of Free Vibration, and Normal Coordinates

The general oscillatory solution is satisfied for a set of n frequencies ω_k . Thus, in order to form a complete solution, we need to take a superposition of all the allowed frequencies.

$$x_i = C_k a_{ik} \exp(-i\omega_k t) \quad (6.35)$$

If we want to trigger only a single frequency, we must find the normal modes. In vector notation,

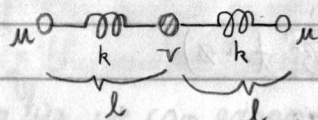
$$L = \frac{1}{2} (\dot{\vec{x}}, \vec{M} \dot{\vec{x}}) - \frac{1}{2} (\vec{x}, \vec{K} \vec{x}) \quad (\text{see 4.44})$$

$$\ddot{\vec{x}} + \Lambda \vec{x} = 0 \quad \Lambda = \vec{M}^{-1} \vec{K} \quad (\text{see 4.45})$$

The normal frequencies are the square roots of the eigenvalues, and the corresponding eigenvectors are the initial condition.

Section 4. Free Vibrations of a Linear Triatomic Molecule (drawn from Jose)

We can treat a water molecule as three particles in a line, connected by equal springs as shown below



$$L = \frac{1}{2} \mu (\dot{x}_1^2 + \dot{x}_3^2) + \frac{1}{2} r \dot{x}_2^2 - \frac{1}{2} k (x_2 - x_1)^2 - \frac{1}{2} k (x_3 - x_2)^2$$

$$= \frac{1}{2} (\dot{\vec{x}}, \vec{M} \dot{\vec{x}}) - \frac{1}{2} (\vec{x}, \vec{K} \vec{x}) \quad (6.1)$$

$$\vec{M} = \begin{bmatrix} \mu & & \\ & r & \\ & & \mu \end{bmatrix} \quad \vec{K} = k \begin{bmatrix} 1 & -1 & \\ -1 & 2 & -1 \\ & -1 & 1 \end{bmatrix} \quad (6.1)$$

i.e.

$$L = \frac{1}{2} \dot{\vec{x}} \vec{M} \dot{\vec{x}} - \frac{1}{2} \vec{x} \vec{K} \vec{x} \quad (6.6)$$

$$\vec{\Lambda} = \vec{K} \vec{M}^{-1} = k \begin{bmatrix} 1/\mu & -1/\mu & \\ -1/r & 2/r & -1/r \\ & -1/\mu & 1/\mu \end{bmatrix} \quad (6.7)$$

$$\det(\Lambda - \lambda I) = \det \begin{bmatrix} 1/\mu - \lambda & -1/\mu & \\ -1/r & 2/r - \lambda & -1/r \\ & -1/\mu & 1/\mu - \lambda \end{bmatrix}$$

$$= (1/\mu - \lambda) [(2/r - \lambda)(1/\mu - \lambda) - 1/\mu r] + 1/\mu (-1/r)(1/\mu - \lambda)$$

$$= (1/\mu - \lambda) [(2/r - \lambda)(1/\mu - \lambda) - 2/\mu r]$$

$$= (1/\mu - \lambda) (\lambda^2 - 2/r - 1/\mu) = 0$$

$$= (1/\mu - \lambda) \cdot \lambda (\lambda - (2/r + 1/\mu)) = 0$$

$$\omega_1 = \sqrt{1/\mu}$$

$$\omega_2 = \sqrt{k(2/r + 1/\mu)}$$

$$\omega_3 = 0$$

$$|\sqrt{1/\mu}\rangle: \begin{bmatrix} 0 & 1/\mu & 0 \\ -1/r & 2/r - 1/\mu & -1/r \\ 0 & -1/\mu & 0 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$b = 0$$

$$a = -c$$

$$|\sqrt{1/\mu}\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 0 \\ -1 \end{bmatrix}$$

$$|\sqrt{k(2/r + 1/\mu)}\rangle: \begin{bmatrix} -2/r & -1/\mu \\ -1/r & 1/\mu - 1/r \\ & -1/\mu & -2/r \end{bmatrix} \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$-2a/r - b/\mu = 0$$

$$-b/\mu - 2c/r = 0$$

$$|\sqrt{k(2/r + 1/\mu)}\rangle = n \begin{bmatrix} 1 \\ -2\mu/r \\ 1 \end{bmatrix}$$

$$|0\rangle: \begin{bmatrix} 1/\mu & -1/\mu \\ -1/r & 2/r & -1/r \\ & -1/\mu & 1/\mu \end{bmatrix} \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$|0\rangle = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

$$\vec{x}(t) = |\sqrt{1/\mu}\rangle [\alpha_1 \exp(i\omega_1 t) + \alpha_1^* \exp(-i\omega_1 t)] + |\sqrt{k(2/r + 1/\mu)}\rangle [\alpha_2 \exp(i\omega_2 t) + \alpha_2^* \exp(-i\omega_2 t)] + |0\rangle [\alpha_3 \exp(i\omega_3 t) + \alpha_3^* \exp(-i\omega_3 t)]$$

Section 5 Forced Vibrations and the Effect of Dissipative Forces

Derivation

$$1. \dot{y}_1 = -\dot{x}_1$$

$$\dot{y}_2 = \dot{x}_2$$

$$L = \frac{\mu}{2} (\dot{y}_1^2 + \dot{y}_2^2) - \frac{k}{2} (y_1^2 + y_2^2)$$

$$M = \begin{bmatrix} \mu & 0 \\ 0 & \mu \end{bmatrix} \quad K = \begin{bmatrix} k & 0 \\ 0 & k \end{bmatrix}$$

$$L = \frac{1}{2} \dot{\vec{y}}^T M \dot{\vec{y}} - \frac{1}{2} \vec{y}^T K \vec{y}$$

$$\Lambda = M^{-1}K = \begin{bmatrix} 1/\mu & 0 \\ 0 & 1/\mu \end{bmatrix} \begin{bmatrix} k & 0 \\ 0 & k \end{bmatrix} = \begin{bmatrix} k/\mu & 0 \\ 0 & k/\mu \end{bmatrix}$$

$$\det \Lambda = (k/\mu - \lambda)^2 = 0$$

$$\lambda = k/\mu$$

$$\omega_1 = \sqrt{k/\mu}$$

$$|\sqrt{k/\mu}\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

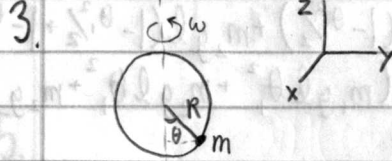
$$|\sqrt{k/\mu}'\rangle = \begin{bmatrix} 0 \\ -1 \end{bmatrix}$$

Which corresponds to moving one of the outside particles

$$\vec{x}(t) = |\sqrt{k/\mu}\rangle (\alpha_1 \exp(i\omega_1 t) + \alpha_1^* \exp(-i\omega_1 t))$$

$$+ |\sqrt{k/\mu}'\rangle (\alpha_2 \exp(i\omega_1 t) + \alpha_2^* \exp(-i\omega_1 t))$$

Problems



$$a. x = -R \sin(\omega t) \sin \theta$$

$$y = R \cos(\omega t) \sin \theta$$

$$z = -R \cos \theta$$

$$L = \frac{m}{2} (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - mgz$$

$$= \frac{m}{2} (R^2 \omega^2 \sin^2 \theta + R^2 \dot{\theta}^2) + mgR \cos \theta$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{d}{dt} (mR^2 \dot{\theta}) - (mR^2 \omega^2 \sin \theta \cos \theta - mgR \sin \theta) = 0$$

$$\dot{\theta} - \omega^2 \sin \theta \cos \theta + \frac{g}{R} \sin \theta = 0$$

b. At equilibrium, $\dot{\theta} = \ddot{\theta} = 0$

$$\omega^2 \cos \theta = \frac{g}{R}$$

$$\omega_c = \sqrt{g/R}$$

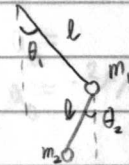
since the bottom of the hoop is the equilibrium position ($\theta = 0$)

c. $\omega > \sqrt{g/R}$

$$\omega^2 = \omega_c^2 / \cos \theta$$

$$\theta_{eq} = \cos^{-1}(\omega_c / \omega)$$

4.



$$x_1 = l \sin \theta_1$$

$$x_2 = l \sin \theta_1 + l \sin \theta_2$$

$$\dot{x}_1 = l \dot{\theta}_1 \cos \theta_1$$

$$\dot{x}_2 = l \dot{\theta}_1 \cos \theta_1 + l \dot{\theta}_2 \cos \theta_2$$

$$y_1 = l \cos \theta_1$$

$$y_2 = l \cos \theta_1 + l \cos \theta_2$$

$$\dot{y}_1 = -l \dot{\theta}_1 \sin \theta_1$$

$$\dot{y}_2 = -l \dot{\theta}_1 \sin \theta_1 - l \dot{\theta}_2 \sin \theta_2$$

$$L = \frac{m_1}{2} (l^2 \dot{\theta}_1^2) + \frac{m_2}{2} (l^2 \dot{\theta}_1^2 + 2l^2 \dot{\theta}_1 \dot{\theta}_2 \cos(\theta_1 - \theta_2) + l^2 \dot{\theta}_2^2)$$

$$- m_1 g l \cos \theta_1 - m_2 g (l \cos \theta_1 + l \cos \theta_2)$$

Working in small angle approximation

$$L \approx \frac{m_1}{2} (l^2 \dot{\theta}_1^2) + \frac{m_2}{2} (l^2 \dot{\theta}_1^2 + l^2 \dot{\theta}_2^2) + m_1 g l (1 - \theta_1^2/2) + m_2 g l (1 - \theta_1^2/2 + 1 - \theta_2^2/2)$$

$$\approx \frac{1}{2} (m_1 \dot{\theta}_1^2 + m_2 \dot{\theta}_1^2 + m_2 \dot{\theta}_2^2 - 2m_2 \dot{\theta}_1 \dot{\theta}_2) - \frac{1}{2} (m_1 g l \theta_1^2 + m_2 g l \theta_1^2 + m_2 g l \theta_2^2)$$

$$M = \begin{bmatrix} l^2(m_1+m_2) & -2m_2 l^2 \\ -2m_2 l^2 & l^2 m_2 \end{bmatrix}$$

$$K = \begin{bmatrix} g l (m_1+m_2) & \\ & g l m_2 \end{bmatrix}$$

$$L \approx \frac{1}{2} \dot{\vec{x}}^T M \dot{\vec{x}} - \frac{1}{2} \vec{x}^T K \vec{x}$$

$$\det(\Lambda - \lambda I) = \det(K - \lambda M) = 0$$

$$\det \begin{bmatrix} g l (m_1+m_2) - l^2 \lambda (m_1+m_2) & m_2 l^2 \lambda \\ m_2 l^2 \lambda & g l m_2 - l^2 \lambda m_2 \end{bmatrix} = 0$$

$$= (m_1+m_2)(g l - l^2 \lambda) m_2 (g l - l^2 \lambda) - m_2^2 l^4 \lambda^2 = 0$$

$$\lambda = \frac{g(m_1+m_2)}{l m_1} \pm \frac{g}{l} \sqrt{\frac{m_1 m_2 + m_2^2}{m_1^2}}$$

$$= \omega_0^2 \left(\frac{M}{m_1} \pm \sqrt{\frac{M m_2}{m_1^2}} \right)$$

$$\omega_0^2 = g/l$$

$$M = m_1 + m_2$$

$$| \omega_0 \left(\frac{M}{m_1} \pm \sqrt{\frac{M m_2}{m_1^2}} \right)^{1/2} | \begin{bmatrix} g l M - \omega^2 l^2 M & \omega^2 m_2 l^2 \\ \omega^2 m_2 l^2 & g l m_2 - \omega^2 l^2 m_2 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$a(g l M - \omega^2 l^2 M) + b \omega^2 m_2 l^2 = 0$$

$$a \omega^2 m_2 l^2 + b(g l m_2 - \omega^2 l^2 m_2) = 0$$

$$a = -\frac{b m_2 (g l - \omega^2 l^2)}{m_2 \omega^2 l^2}$$

$$\frac{M(g l - \omega^2 l^2)^2}{\omega^2 l^2} + b \omega^2 m_2 l^2 = 0$$

$$b = -\frac{(g l - \omega^2 l^2)^2}{\omega^4 l^4}$$

$$| \omega_+ \rangle = \begin{bmatrix} \sqrt{m_2/M} \\ 1 \end{bmatrix}$$

$$| \omega_- \rangle = \begin{bmatrix} \sqrt{m_2/M} \\ -1 \end{bmatrix}$$

5.

a. The general solution is

$$\begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} (\alpha_1 \exp(i\omega_1 t) + \alpha_1^* \exp(-i\omega_1 t))$$

$$+ \begin{pmatrix} 1 \\ -2\mu/r \\ 1 \end{pmatrix} (\alpha_2 \exp(i\omega_2 t) + \alpha_2^* \exp(-i\omega_2 t))$$

$$+ \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} (a + vt)$$

$$\begin{pmatrix} 0 \\ a_0 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} (\alpha_1 + \alpha_1^*) + \begin{pmatrix} 1 \\ -2\mu/r \\ 1 \end{pmatrix} (\alpha_2 + \alpha_2^*) + \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} a$$

$$\begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} (i\omega_1 \alpha_1 - i\omega_1 \alpha_1^*) + \begin{pmatrix} 1 \\ -2\mu/r \\ 1 \end{pmatrix} (i\omega_2 \alpha_2 - i\omega_2 \alpha_2^*) + \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} v$$

$$0 = \alpha_1 + \alpha_1^* + \alpha_2 + \alpha_2^* + a$$

$$0 = i\omega_1 \alpha_1 - i\omega_1 \alpha_1^* + i\omega_2 \alpha_2 - i\omega_2 \alpha_2^* + v$$

$$0 = -\alpha_1 - \alpha_1^* + \alpha_2 + \alpha_2^* + a$$

$$0 = -i\omega_1 \alpha_1 + i\omega_1 \alpha_1^* + i\omega_2 \alpha_2 - i\omega_2 \alpha_2^* + v$$

$$a_0 = -2\mu/r (\alpha_2 + \alpha_2^*) + a$$

$$0 = -2\mu/r (i\omega_2 \alpha_2 - i\omega_2 \alpha_2^*) + v$$

$$a_0 = -2\mu/r (\alpha_2 + \alpha_2^*) - (\alpha_2 + \alpha_2^*)$$

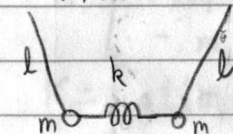
$$0 = -2\mu/r (i\omega_2 \alpha_2 - i\omega_2 \alpha_2^*) - (i\omega_2 \alpha_2 + i\omega_2 \alpha_2^*)$$

$$0 = (-2\mu/r - 1)(i\omega_2)(\alpha_2 - \alpha_2^*) = 0$$

$$\begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} (\exp(i\omega_1 t) - \exp(-i\omega_1 t)) + \begin{pmatrix} 1 \\ -2\mu/r \\ 1 \end{pmatrix} (\exp(i\omega_2 t) - \exp(-i\omega_2 t))$$

$$+ \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} (a_0)$$

$$b. \vec{x} = \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} (\exp(i\omega_1 t) - \exp(-i\omega_1 t)) + \begin{pmatrix} -1/2 \\ 1/2 \\ 1 \end{pmatrix} (\exp(i\omega_2 t) - \exp(-i\omega_2 t)) + \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} v_0 t$$

6.  $m \frac{d^2 x}{dt^2} + kx + mg \sin \theta = 0$

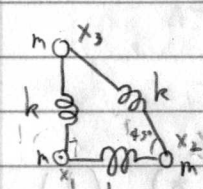
$$m \frac{d^2 x}{dt^2} + kx + mg \frac{x}{l} = 0$$

$$m \ddot{x} = -\left(\frac{k}{m} + \frac{g}{l}\right)x$$

$$\omega = \sqrt{\frac{k}{m} + \frac{g}{l}}$$

$$T = \frac{2\pi}{\omega} = \frac{2\pi}{\sqrt{\frac{3}{.23} + \frac{10}{.8}}} = 1.24 \text{ s}$$

7.

8.  $(n_1, \epsilon_1) = (x_1, y_1)$
 $(n_2, \epsilon_2) = (x_2 - a, y_2)$
 $(n_3, \epsilon_3) = (x_3, y_3 - a)$

$$T = \frac{m}{2} (\dot{n}_1^2 + \dot{\epsilon}_1^2 + \dot{n}_2^2 + \dot{\epsilon}_2^2 + \dot{n}_3^2 + \dot{\epsilon}_3^2)$$

$$V_{12} = \frac{k}{2} [\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2} - a]^2$$

$$= \frac{k}{2} [\sqrt{(n_2 + a - n_1)^2 + (\epsilon_2 - \epsilon_1)^2} - a]^2$$

$$\approx \frac{k}{2} [a^2 (1 + \frac{n_2 - n_1}{a})^2 - a]^2$$

$$\approx \frac{k}{2} [a (1 + \frac{n_2 - n_1}{a}) - a]^2 = \frac{k}{2} (n_2 - n_1)^2$$

$$V_{13} = \frac{k}{2} (\epsilon_3 - \epsilon_1)^2$$

$$V_{23} = \frac{k}{2} [\sqrt{(x_3 - x_2)^2 + (y_3 - y_2)^2} - \sqrt{2}a]^2$$

$$= \frac{k}{2} [(\frac{n_3 - n_2 - a}{a})^2 + (\frac{\epsilon_3 + a - \epsilon_2}{a})^2 - \sqrt{2}]^2$$

$$\approx \frac{k}{2} [a^2 (1 - \frac{n_3 - n_2}{a})^2 + a^2 (1 - \frac{\epsilon_3 - \epsilon_2}{a})^2 - \sqrt{2}a]^2$$

$$= \frac{k}{4} [(n_3 - n_2)^2 - (\epsilon_3 - \epsilon_2)^2]^2$$

$$V = \frac{k}{2} [(n_2 - n_1)^2 + (\epsilon_3 - \epsilon_1)^2 + \frac{1}{2} (n_3 - n_2 - \epsilon_3 + \epsilon_2)^2]$$

$$V - \lambda T = \begin{bmatrix} k - m\lambda & 0 & -k & 0 & 0 & 0 \\ 0 & k - m\lambda & 0 & 0 & 0 & -k \\ -k & 0 & \frac{3}{2}m\lambda & -\frac{k}{2} & -\frac{k}{2} & \frac{k}{2} \\ 0 & 0 & -\frac{k}{2} & \frac{1}{2}m\lambda & \frac{k}{2} & -\frac{k}{2} \\ 0 & 0 & -\frac{k}{2} & \frac{k}{2} & \frac{1}{2}m\lambda & -\frac{k}{2} \\ 0 & -k & \frac{k}{2} & -\frac{k}{2} & -\frac{k}{2} & \frac{3}{2}m\lambda \end{bmatrix}$$

9. $\Delta E = 5 \text{ J}$

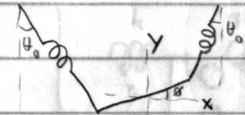
$$\Delta E = \frac{1}{2} kx^2$$

$$k = 10 \text{ N/m}$$

10.

a.

11.



$$T = \frac{m}{2} (\dot{x}^2 + \dot{y}^2) + \frac{1}{2} I \dot{\theta}^2$$

$$I = \int_{-l/2}^{l/2} x^2 \rho dx = \frac{m}{l} \cdot \frac{1}{3} x^3 \Big|_{-l/2}^{l/2} = \frac{m}{3l} \cdot \frac{2l^3}{8} = \frac{ml^2}{12}$$

$$T = \frac{m}{2} (\dot{x}^2 + \dot{y}^2) + \frac{ml^2 \dot{\theta}^2}{24}$$

$$V = \frac{k}{2} [(x + l \cos \theta_0 - x - l \cos \theta_0)^2 + (y + l \sin \theta_0 - y - l \sin \theta_0)^2 - b]^2$$

$$V = \frac{k}{2} [\sqrt{(l/2 + b \sin \theta_0 - x - l \cos \theta_0)^2 + (b \cos \theta_0 - y - l \sin \theta_0)^2} - b]^2$$

$$\approx \frac{k}{2} [((b \sin \theta_0 - x)^2 + (b \cos \theta_0 - y - l \sin \theta_0)^2)^{1/2} - b]^2$$

Keeping linear terms in x, y, θ

$$\approx \frac{k}{2} [x \sin \theta_0 + (y + l \sin \theta_0) \cos \theta_0]^2$$

$$+ \frac{k}{2} [x \sin \theta_0 - (y - l \sin \theta_0) \cos \theta_0]^2$$

$$T = \begin{bmatrix} m & & \\ & m & \\ & & m/3 \end{bmatrix}$$

$$V = \begin{bmatrix} \sin^2 \theta_0 & 0 & -\sin \theta_0 \cos \theta_0 \\ 0 & \cos^2 \theta_0 & 0 \\ -\sin \theta_0 \cos \theta_0 & 0 & \cos^2 \theta_0 \end{bmatrix} \cdot 2k$$

$$\det(V - \lambda T) = \begin{vmatrix} 2k \sin^2 \theta_0 - \lambda m & 0 & -2k \sin \theta_0 \cos \theta_0 \\ 0 & 2k \cos^2 \theta_0 - \lambda m & 0 \\ -2k \sin \theta_0 \cos \theta_0 & 0 & 2k \cos^2 \theta_0 - \lambda m/3 \end{vmatrix}$$

$$= (2k \sin^2 \theta_0 - \lambda m)(2k \cos^2 \theta_0 - \lambda m)(2k \cos^2 \theta_0 - \lambda m/3)$$

$$- (2k \sin \theta_0 \cos \theta_0)^2 (2k \cos^2 \theta_0 - \lambda m)$$

$$= (2k \cos^2 \theta_0 - \lambda m) [4k^2 \sin^2 \theta_0 \cos^2 \theta_0 - 2km \lambda \cos^2 \theta_0 + m^2/3 - 4k^2 \sin^2 \theta_0 \cos^2 \theta_0]$$

$$= (2k \cos^2 \theta_0 - \lambda m) \cdot m^2/3 (m \lambda - 2k \sin^2 \theta_0 - 6k \cos^2 \theta_0)$$

$$= (2k \cos^2 \theta_0 - \lambda m) \cdot m^2/3 (m \lambda - 2k - 4k \cos^2 \theta_0)$$

$$\lambda = 0, \quad \frac{2k}{m} \cos^2 \theta_0, \quad \frac{2k}{m} (1 + 2 \cos^2 \theta_0)$$

$$10 \rangle: \begin{bmatrix} 2k \sin^2 \theta_0 & 0 & -2k \sin \theta_0 \cos \theta_0 \\ 0 & 2k \cos^2 \theta_0 & 0 \\ -2k \sin \theta_0 \cos \theta_0 & 0 & 2k \cos^2 \theta_0 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$b = 0$$

$$2ka \sin^2 \theta_0 - 2kc \sin \theta_0 \cos \theta_0 = 0$$

$$a \sin \theta_0 = c \cos \theta_0$$

$$10 \rangle = \begin{bmatrix} \cos \theta_0 \\ 0 \\ -\sin \theta_0 \end{bmatrix}$$

$$| \frac{2k}{m} \cos^2 \theta_0 \rangle: \begin{bmatrix} 2k \sin^2 \theta_0 - 2k \cos^2 \theta_0 & 0 & -2k \sin \theta_0 \cos \theta_0 \\ 0 & 0 & 0 \\ -2k \sin \theta_0 \cos \theta_0 & 0 & 4k/3 \cos^2 \theta_0 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$| \frac{2k}{m} \cos^2 \theta_0 \rangle = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$$

$$| \frac{2k}{m} (1 + 2 \cos^2 \theta_0) \rangle: \begin{bmatrix} 2k \sin^2 \theta_0 - 2k - 4k \cos^2 \theta_0 & 0 & -2k \sin \theta_0 \cos \theta_0 \\ 0 & -2k - 2k \cos^2 \theta_0 & 0 \\ -2k \sin \theta_0 \cos \theta_0 & 0 & 4k/3 \cos^2 \theta_0 - 2k/3 \end{bmatrix}$$

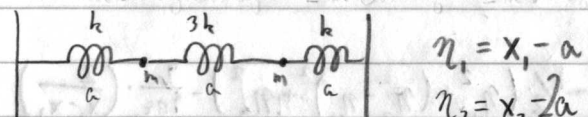
$$a(\sin^2 \theta_0 - 1 - 2 \cos^2 \theta_0) - c \sin \theta_0 \cos \theta_0 = 0$$

$$-3a \cos^2 \theta_0 = c \sin \theta_0 \cos \theta_0$$

$$c \sin \theta_0 = -2c \sin \theta_0 = -3a \cos \theta_0$$

$$| \frac{2k}{m} (1 + 2 \cos^2 \theta_0) \rangle = \begin{bmatrix} \sin \theta_0 \\ 0 \\ -3 \cos \theta_0 \end{bmatrix}$$

12.



$$\eta_1 = x_1 - a$$

$$\eta_2 = x_2 - 2a$$

$$T = \frac{m}{2} (\dot{x}_1^2 + \dot{x}_2^2) = \frac{m}{2} (\dot{\eta}_1^2 + \dot{\eta}_2^2)$$

$$V = \frac{k}{2} (x_1 - a)^2 + \frac{3k}{2} (x_2 - x_1 - a)^2 + \frac{k}{2} (x_2 - 2a)^2$$

$$= \frac{k}{2} (\eta_1^2) + \frac{3k}{2} (\eta_2 - \eta_1)^2 + \frac{k}{2} (\eta_2^2)$$

$$= \frac{k}{2} [\eta_1^2 + 3(\eta_2^2 - 2\eta_1 \eta_2 + \eta_1^2) + \eta_2^2]$$

$$= \frac{k}{2} [4\eta_1^2 - 3\eta_1 \eta_2 - 3\eta_2 \eta_1 + 4\eta_2^2]$$

$$T = \begin{bmatrix} m & \\ & m \end{bmatrix}$$

$$V = \begin{bmatrix} 4k & -3k \\ -3k & 4k \end{bmatrix}$$

$$\mathcal{L} = \frac{1}{2} \dot{\eta} T \dot{\eta} + \frac{1}{2} \eta \cdot V \cdot \eta$$

$$\det(V - \lambda T) = \det \begin{bmatrix} 4k - \lambda m & -3k \\ -3k & 4k - \lambda m \end{bmatrix}$$

$$= (4k - \lambda m)^2 - 9k^2 = \lambda^2 m^2 - 8k\lambda m + 7k^2$$

$$= (\lambda m - 7k)(\lambda m - k)$$

$$\lambda = -7k/m, k/m$$

$$|-7k/m\rangle: \begin{bmatrix} -3k & -3k \\ -3k & -3k \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$|-7k/m\rangle = \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

$$|k/m\rangle: \begin{bmatrix} 3k & -3k \\ -3k & 3k \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$|k/m\rangle = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

13. $\eta_1 = x_1 - a$

$\eta_2 = x_2 - a$

$$\mathcal{L} = \frac{m}{2}(\dot{\eta}_1^2 + \dot{\eta}_2^2) - \frac{k}{2}(\eta_1^2 + (\eta_2 - \eta_1)^2 + \eta_2^2) - \frac{1}{4\pi\epsilon_0} \left(\frac{2}{x_2 - x_1} \right)$$

16. $y = ax^4$

$$\dot{y} = 4ax^3 \dot{x}$$

$$\mathcal{L} = \frac{m}{2}(\dot{x}^2 + \dot{y}^2) - mgy$$

$$= \frac{m}{2}(\dot{x}^2 + 16a^2 x^6 \dot{x}^2) - mga x^4$$

$$x = x_0 + \eta$$

$$y = a(x_0 + \eta)^4$$

$$\dot{x} = \dot{\eta}$$

$$\dot{y} = 4a(x_0 + \eta)^3 \dot{\eta}$$

$$\mathcal{L} = \frac{m}{2}(\dot{\eta}^2 + 16a^2(x_0 + \eta)^6 \dot{\eta}^2) - mga(x_0 + \eta)^4$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\eta}} \right) - \frac{\partial \mathcal{L}}{\partial \eta} = 0$$

$$\frac{d}{dt} (m\dot{\eta} + 16ma^2(x_0 + \eta)^6 \dot{\eta}) - (3m \cdot 16a^2(x_0 + \eta)^5 \dot{\eta}^2 - 4mga(x_0 + \eta)^3) = 0$$

$$m\ddot{\eta} + 16ma^2(x_0 + \eta)^5 \dot{\eta}^2 + 16ma^2(x_0 + \eta)^6 \ddot{\eta} - 48ma^2(x_0 + \eta)^5 \dot{\eta}^2 + 4mga(x_0 + \eta)^3 = 0$$

$$\ddot{\eta} (1 + 16a^2(x_0 + \eta)^6) + 48a^2(x_0 + \eta)^5 \dot{\eta}^2 + 4ga(x_0 + \eta)^3 = 0$$

Cyclic Coordinates and Conservation Laws

Remember, a cyclic coordinate is one that does not appear in the Lagrangian, meaning the coordinate is ignorable. We can convince ourselves that cyclic coordinates are conserved from the Hamiltonian.

$$\frac{\partial H}{\partial q_i} = 0 \implies \dot{p}_i = 0 \implies p_i = \text{const}$$

without the use of Hamilton's equations (8.8)